

Manifestations of top compositeness at colliders

This article has been downloaded from IOPscience. Please scroll down to see the full text article.

JHEP05(2009)022

(<http://iopscience.iop.org/1126-6708/2009/05/022>)

[The Table of Contents](#) and [more related content](#) is available

Download details:

IP Address: 80.92.225.132

The article was downloaded on 03/04/2010 at 09:20

Please note that [terms and conditions apply](#).

Manifestations of top compositeness at colliders

Kunal Kumar,^{a,b} Tim M.P. Tait^{a,b} and Roberto Vega-Morales^{a,b}

^a*Northwestern University,*

2145 Sheridan Road, Evanston, IL 60208, U.S.A.

^b*HEP Division, Argonne National Laboratory,*

Argonne, IL 60439, U.S.A.

E-mail: kunal.kunalkumar@gmail.com, tait@anl.gov,

rvegamorales@gmail.com

ABSTRACT: We explore the possibility that the right-handed top quark is composite, identifying possible signatures of compositeness and how they might manifest themselves at the LHC and Tevatron. We perform a complete analysis of the dimension six modifications of the top coupling to gluons and find that cancellations among operators in the $t\bar{t}$ rate allow for very low compositeness scales, but this can be drastically improved by looking at kinematic distributions. Turning to the LHC, we examine four top production from a dimension six four-top operator and estimate the LHC with 100 fb^{-1} collected luminosity to be sensitive to compositeness scales as high as 5 TeV.

KEYWORDS: Beyond Standard Model, Heavy Quark Physics, Hadronic Colliders, Technicolor and Composite Models

ARXIV EPRINT: [0901.3808](https://arxiv.org/abs/0901.3808)

Contents

1	Introduction	1
2	Top operators describing new physics	2
3	Top pairs at the Tevatron	3
4	Four tops at the LHC	6
5	Outlook	8

1 Introduction

With the advent of the Large Hadron Collider era, it is important to understand in detail how physics beyond the Standard Model might manifest itself in LHC data. While detailed exploration of the impact of models such as supersymmetric ones is well underway, it is also important to be ready for surprises. The LHC energies are sufficiently beyond any we have previously explored directly that there is room for new phenomena that may not be directly connected to outstanding theoretical problems such as the stability of the electroweak scale. In this article, we explore the idea that the top quark is composite, composed of some new constituent particles (“preons”) bound together by a new confining force [1–3]. As the most recently discovered ingredient of the Standard Model, the top is the quark about which we currently have the least information. This opens the door to the possibility of large new effects in its interactions being relatively weakly constrained by existing measurements, and its large mass may be a clue that it is different from other Standard Model fermions. In fact, many models exist in which the top quark is composite [4] in order to motivate its large mass, and lead to a variety of interesting signals at the LHC [5].

We choose to work generically in a framework in which the right-handed top is composite, without getting attached to any specific model. Our hope is to identify interesting phenomena and features which are not specifically linked to any particular model, but might reasonably be expected to occur in a broad class of models in which the top is composite. We use effective field theory as a framework with which we can explore the most important operators which can modify the top’s interactions with itself and with other particles of the Standard Model. As a result, the bounds we find can in principle be applied to any theory which modifies the top’s interactions, and are not limited to a specific theory of top compositeness, or indeed a theory of top compositeness at all.

This article is organized as follows. In section 2, we present the most important dimension six operators leading to modifications of top quark interactions. They are expected to lead to the dominant effects at “low” energies, below the scale of the new physics. In

section 3, we revisit their effects on $t\bar{t}$ production at the Tevatron. This extends previous results [2], which strictly assumed coefficients as dictated by naive dimensional analysis (NDA) [6] to a general effective theory framework, which flexibly describes the effects of heavy new physics on $t\bar{t}$ production. As a result, it can describe theories of top compositeness which deviate from NDA expectations, or indeed *any* theory which predicts a modification of the top's coupling to gluons. We also examine kinematic distributions and find that they are much more powerful than the rate itself in revealing the presence of new physics. In section 4 we examine the possibility that the operators modifying the properties of top lie at some what higher energy scales, and deduce their contribution to the production rate of four top quarks at the LHC. We conclude in section 5.

2 Top operators describing new physics

We parameterize the modifications of top interactions through heavy new physics in the form of higher dimensional operators which we add to the Standard Model action [7]. This is justified based on the fact that top observables to date do not show any radical deviation from the SM predictions and it is reasonable to expect that the scale of compositeness is probably somewhat higher than the typical energies probed. Because at the current time and in the near future our information about top comes exclusively from hadron colliders, we are most interested in operators capturing modifications to the interactions of the top quark with itself and/or light quarks and gluons, since these lead to the largest effects at a hadron machine. The relevant terms in the effective Lagrangian are,

$$\mathcal{L} = \mathcal{L}_{\text{SM}} + \frac{1}{\Lambda^2} \{ \mathcal{O}_t + \mathcal{O}_1 + \mathcal{O}_2 + \mathcal{O}_3 + \dots \} + \mathcal{O} \left(\frac{1}{\Lambda^4} \right) \quad (2.1)$$

where we define each operator \mathcal{O}_i below, and additional operators with mass dimension > 6 will have suppressed contributions to the processes we consider, which take place at energies $\lesssim \Lambda$.

If t_R and no other Standard Model component is composite, NDA argues that the most significant operator will involve four top quarks [1],

$$\mathcal{O}_t = g^2 [\bar{t}^i \gamma^\mu P_R t_j] [\bar{t}^k \gamma_\mu P_R t_l] \quad (2.2)$$

where γ^μ are the Dirac gamma matrices, P_R is the right-chiral projector, and g^2/Λ^2 is the coupling of this new interaction. In the compositeness picture, it can be understood that g/Λ represents the amplitude to create the composite field, and Λ itself characterizes the energy scale at which further elements of the composite sector become important. The effective theory is sensible at energies below Λ provided $g \lesssim 4\pi$. We consider two options for gauge-invariant contractions of the color indices: contractions of two octets $(T^a)_i^j (T^a)_k^l$ and contractions of two singlets $\delta_i^j \delta_k^l$. Note that operators constructed by replacing one or more of the top quarks with a right-handed up or charm quark are also possible, and could lead to more stringent bounds from flavor-violating processes. For simplicity, we assume such operators have small enough coefficients that they are not important.

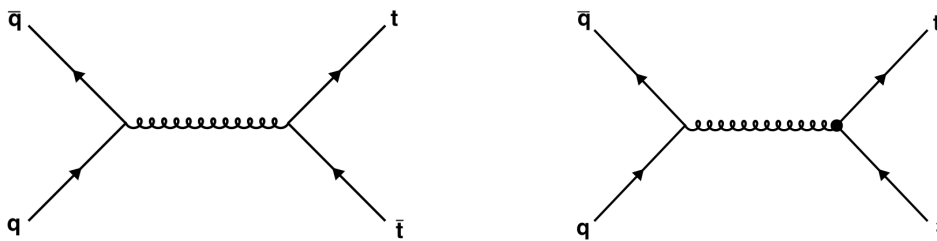


Figure 1. Feynman diagrams for the process $q\bar{q} \rightarrow t\bar{t}$. The shaded blob in the diagram on the right represents the insertion of the new physics g - t - \bar{t} vertex.

While NDA argues that the four top operator is likely to be the most important one in a theory in which the top is composite, it is very difficult to constrain directly at the Tevatron, where the energy requirement to make four real top quarks renders the rate negligibly small. Thus, we turn to sub-leading operators which can contribute to observable processes, such as $t\bar{t}$ production. These operators represent the first potential glimpse that the top is not point-like, as probed by short wavelength fundamental particles such as gluons. The three dimension six operators are [8, 9],

$$\begin{aligned}\mathcal{O}_1 &= g_1 g_S [(H\bar{Q}_3) \sigma^{\mu\nu} \lambda^a P_R t] G_{\mu\nu}^a + H.c. \\ \mathcal{O}_2 &= g_2 g_S [\bar{t} \gamma^\mu \lambda^a D^\nu P_R t] G_{\mu\nu}^a + H.c. \\ \mathcal{O}_3 &= g_3 g_S [\bar{t} \gamma^\mu \lambda^a P_R t] \sum_q [\bar{q} \gamma_\mu \lambda^a q]\end{aligned}\tag{2.3}$$

where g_1 , g_2 , and g_3 parameterize the relative sizes of the three operators, whose over-all energy scale is set by Λ , and we have assumed universal coupling between all light quarks and top in \mathcal{O}_3 . NDA provides the estimates $g_1(\Lambda) \sim O(1/g)$ and $g_2(\Lambda), g_3(\Lambda) \sim O(1)$. The operator \mathcal{O}_1 is a chromo-magnetic moment for the top [10], once the Higgs H is replaced with its VEV. \mathcal{O}_2 is induced by the four top operator of Eq (2.2) through renormalization, and it was in that limit (with $g_2(\Lambda) = 0$) that the previous study used its influence on top pair production at the Tevatron to place a limit on top compositeness [2]. The equations of motion relate \mathcal{O}_3 to $\mathcal{O}_2 + \mathcal{O}_2^\dagger$. We can thus consider complex g_1 and g_2 and $g_3 = 0$ without loss of generality, since a non-zero (real) g_3 can be recast through the equations of motion into a shift in the real part of g_2 . For the purposes of our study, we do not restrict ourselves to the values motivated by NDA, allowing for general complex values of g_1 and g_2 .

3 Top pairs at the Tevatron

As alluded above, at the Tevatron the most promising measurements to constrain new physics in the top system come from $t\bar{t}$ production. In this section we derive constraints from the inclusive rate of top pair production, $\sigma(t\bar{t})$. These constraints turn out to be relatively weak because of the possibility of cancellations between the effects of \mathcal{O}_1 and \mathcal{O}_2 , but we show how kinematic distributions can improve this situation, and lead to much stronger bounds.

The inclusive $t\bar{t}$ cross section is measured by CDF [11] and D0 [12] to be,

$$\sigma(t\bar{t})_{\text{CDF}} = 7.0 \pm 0.3 \pm 0.4 \pm 0.4 \text{ pb} \quad \sigma(t\bar{t})_{\text{D0}} = 7.62 \pm 0.85 \text{ pb}, \quad (3.1)$$

(quoted at $m_t = 172.6 \text{ GeV}$) where the errors are (in order) statistical, systematic, and arising from the luminosity measurement. Both are slightly higher than the Standard Model prediction,

$$\sigma(t\bar{t})_{\text{SM}} = 6.6 \pm 0.8 \text{ pb}, \quad (3.2)$$

(we combine results from both references of [13], to obtain this estimate), but not significantly so. The CDF measurement has slightly smaller error bars, and is slightly closer to the SM, and thus results in the stricter bound. In order to be conservative, we base our limit on it, combining the various errors in quadrature to arrive at $\sigma(t\bar{t})_{\text{exp}} = 7.0 \pm 0.61 \text{ pb}$.

At the Tevatron, $q\bar{q}$ initial states provide more than 80% of the rate, and thus the dominant effect of the operators modifying the $g-t\bar{t}$ interaction results from an insertion of the modified vertex in the usual SM production diagram, as shown in figure 1. Since SM production does a reasonably good job describing the data, we work at leading order in the $1/\Lambda^2$, in which the modified vertex interferes with the usual SM production diagram. We neglect the gluon-initiated graph, which at the Tevatron amounts to an error of roughly 10% or so in our estimates. In the composite top theory, this correction looks like a hard gluon interacting directly with some of the colored preons inside the top, probing its internal structure. Up to $O(1/\Lambda^2)$, the partonic cross section $\hat{\sigma}(q\bar{q} \rightarrow t\bar{t})$ is given by,

$$\hat{\sigma} \propto |\mathcal{M}|^2 = |\mathcal{M}_{\text{SM}} + \mathcal{M}_{\text{NP}}|^2 = |\mathcal{M}_{\text{SM}}|^2 + 2\text{Re}[\mathcal{M}_{\text{SM}}^* \mathcal{M}_{\text{NP}}] + O\left(\frac{1}{\Lambda^4}\right), \quad (3.3)$$

with the leading new physics correction arising from the cross term between the standard model and new physics amplitudes. The resulting cross section can be rewritten to be proportional to σ_{SM} ,

$$\hat{\sigma} = \sigma_{\text{SM}} \left(1 + \text{Re} \frac{g_1(16vms^2) + g_2(4m^2s^2 + s^3 + s(s+2t-2m^2)^2)}{2\Lambda^2(2m^4 + s^2 - 4m^2t + 2st + 2t^2)} \right) + O\left(\frac{1}{\Lambda^4}\right). \quad (3.4)$$

where s and t are the Mandelstam variables, and m is the top quark mass. Note that at leading order, only the real parts of g_1 and g_2 contribute.

We simulate the new physics effect by generating SM $t\bar{t}$ events using Madevent [14] and subsequently reweighting them by eq. (3.4) on an event-by-event basis, according to the generated kinematics. We vary the real parts of g_1 and g_2 , with Λ fixed to 500 GeV, and in figure 2 we present the result for the inclusive $t\bar{t}$ cross section in the plane of g_1 and g_2 . The resulting cross section is somewhat more sensitive to \mathcal{O}_2 than to \mathcal{O}_1 . Also shown are the one sigma contours around the central experimental measurement, revealing that if g_1 and g_2 have opposite signs, their effects in $\sigma(t\bar{t})$ may cancel, resulting in no large net effect in the inclusive rate, despite the presence of low scale new physics.

We expect to better constrain g_1 and g_2 by comparing the distributions of the $t\bar{t}$ invariant mass [15] and top rapidity with experimental data. To illustrate the effectiveness

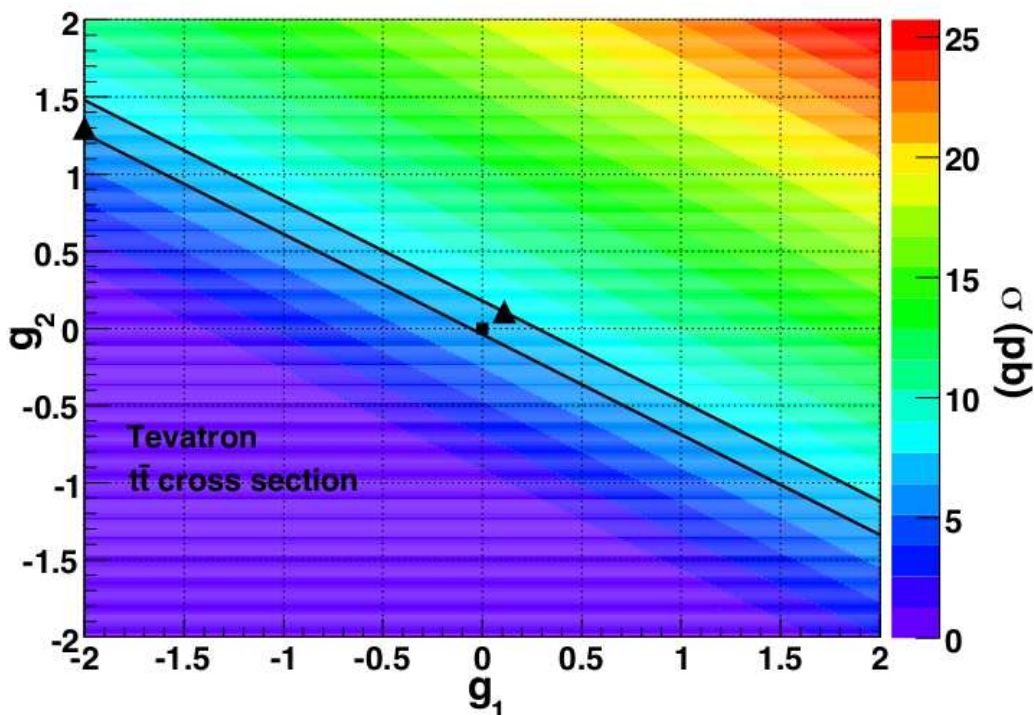


Figure 2. The inclusive top pair production cross section in the plane of g_1 and g_2 . The region between the two black lines is consistent with Tevatron $t\bar{t}$ cross section measurements at 1σ , and the Standard Model prediction is represented by the black square. The black triangles denote two points which lead to a cross section within 1σ of the experimental results, and for which we present detailed distributions in figure 3.

of these distributions in revealing the presence of physics beyond the Standard Model, we consider two points in the parameter space for which we compute the $t\bar{t}$ invariant mass and top rapidity distributions. We choose our two points to lie near the boundary of experimental error, thus leading to the largest deviation from the Standard Model not already in contradiction with the measured $t\bar{t}$ cross section. These points are indicated by black triangles on figure 2,

$$\text{Point 1 : } (g_1, g_2) = (0.11, 0.11) \quad \text{Point 2 : } (g_1, g_2) = (-2.0, 1.3) . \quad (3.5)$$

Point 1 represents a case where both operators do not result in any cancellation in the inclusive rate, whereas Point 2 is a case where moderate cancellations result in a $t\bar{t}$ cross section consistent with the Standard Model.

In figure 3, we present the invariant mass and rapidity distributions for both points. Our results reveal that despite the mild effect on the $t\bar{t}$ rate, the effect of Point 2 on the $t\bar{t}$ invariant mass distribution is striking – in fact, so striking that this point is certainly already ruled out by the searches for $t\bar{t}$ resonances [16]. While that data is not presented in such a way as to allow an outsider to perform the analysis, we encourage the experimental collaborations to consider deriving such bounds as a by-product of the resonance searches.

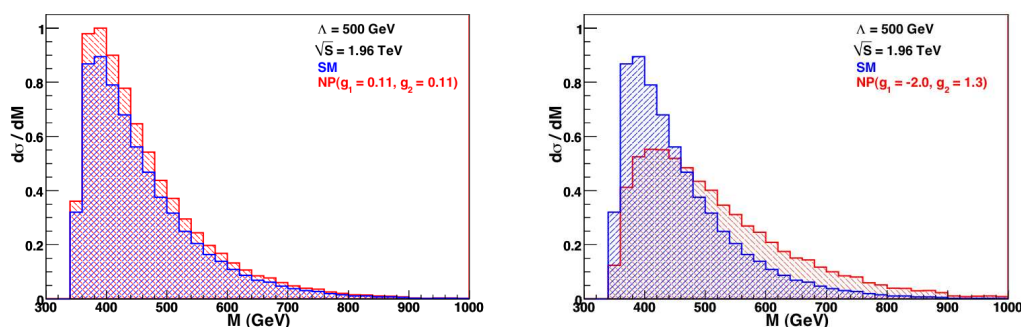


Figure 3. Invariant mass distribution of $t\bar{t}$ for $(g_1 = g_2 = 0.11)$ and $(g_1 = -2.0, g_2 = 1.3)$

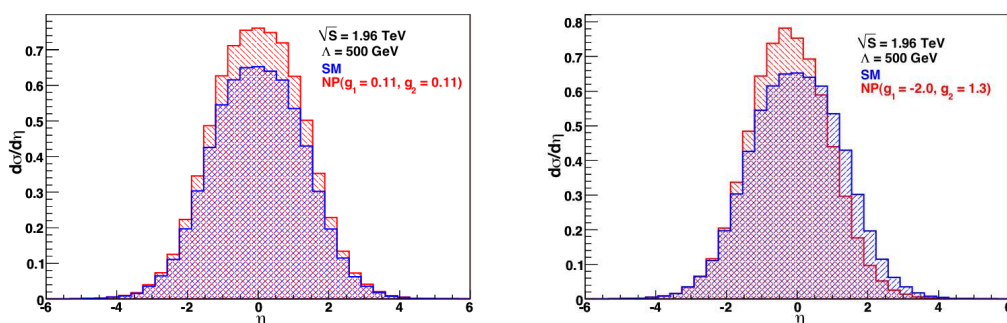


Figure 4. Rapidity distribution of the top quark for $(g_1 = g_2 = 0.11)$ and $(g_1 = -2.0, g_2 = 1.3)$

The effects on the top rapidity distributions are somewhat more modest, but with the interesting feature that because \mathcal{O}_2 distinguishes the top from the anti-top, it can lead an asymmetric top rapidity distribution, and a deviation in the forward-backward asymmetry of the top [17].

4 Four tops at the LHC

At the LHC, energies are sufficiently large to consider processes mediated by the four top operator of eq. (2.2), namely $pp \rightarrow t\bar{t}t\bar{t}$ through processes such as $pp \rightarrow t\bar{t}^*$ followed by $\bar{t}^* \rightarrow \bar{t}t\bar{t}$ (see figure 5 for a representative Feynman diagram). As four top production is characterized by higher energies than $t\bar{t}$, it naturally provides a larger lever-arm with which to study higher dimensional operators. In this section, we consider larger values of Λ , such that the effective theory truncated at order $1/\Lambda^2$ remains an accurate description at LHC energies. Thus, for the purposes of this section we consider the case where the new physics results in effects probably too small to be seen at the Tevatron, and the LHC is the first place where non-standard interactions for top may be visible.

Since the Standard Model prediction for four top production at the LHC is very small (perhaps unobservably so [2, 18]), any positive signal at the LHC will arise in a regime where the new physics amplitudes for $gg \rightarrow t\bar{t}t\bar{t}$ are much larger than the sum of the Standard

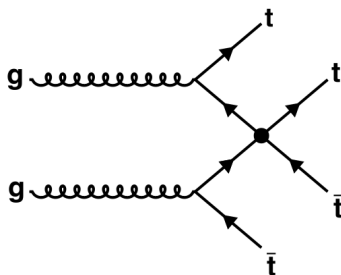


Figure 5. Representative Feynman diagram illustrating the contribution from a four top operator to four top production at the LHC.

Model amplitudes. Thus, for this process to be interesting, we expect that the $O(1/\Lambda^4)$ contribution to the cross section from the square of diagrams such as the one in figure 5 will dominate over both the Standard Model contribution and the $O(1/\Lambda^2)$ and $O(1/\Lambda^4)$ interference terms with dimension 8 operators. We neglect these smaller contributions in presenting our results. A very similar study (for the octet case) was previously presented in [3], including interference with the SM amplitudes, and the results are similar to the ones found here.

We implement the four top operator into MadEvent by introducing an auxiliary field which mimics the dimension six operator at low energies (we verify that the result scales as expected with the mass of the auxiliary field as a cross check). For a theory with a composite top, with strongly coupled new physics, we set $g = 4\pi$, as suggested by NDA; the rates for theories with different values of g are easily obtained by scaling our results by $g^4/(4\pi)^4$. The resulting cross section, as a function of Λ , is shown in figure 6.

The Standard Model prediction for the four top rate is about 3 fb [2], and we see that new physics can lead to enhancements by many orders of magnitude over the SM. A previous study [2] based on like-sign leptonic W decays and 100 fb^{-1} of integrated luminosity concluded that four tops are observable over backgrounds (dominantly di-boson production plus jets and $t\bar{t}$) provided the raw four top production is greater than 45 fb. From figure 6, this occurs for $\Lambda \leq 5 \text{ TeV}$, and is the largest scale of top compositeness to which the LHC (with 100 fb^{-1}) has sensitivity based on a simple like-sign lepton analysis. Since our production is through a higher dimensional operator and favors relatively boosted tops, it would be interesting to see if hadronic top reconstruction through jet shape variables could provide additional channels with which the four top rate could be measured [19].

Many other theories can lead to an excess in the four top production rate, including theories with color octet scalars [18, 20], color octet vectors [21], color sextets [22], and supersymmetric theories with decay chains through third family sfermions [23]. Theories with additional down-type quarks [24] which decay into W and top could also be revealed by the like-sign W search strategy. In each case, there are additional features in the kinematic distributions which can help distinguish which model is responsible for a positive signal. For example, we expect to reconstruct resonances in $t\bar{t}$ pairs for the octets, resonances in tt ($t\bar{t}$) for the sextets, additional sources of missing energy beyond the leptonic W decays

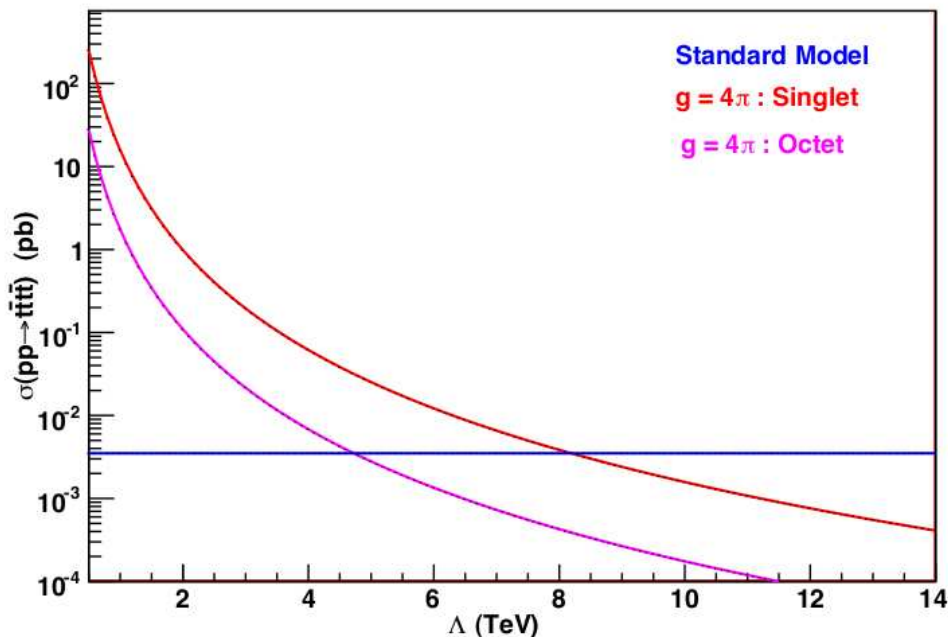


Figure 6. Rate of four top production at the LHC as a function of Λ with $g = 4\pi$. Also indicated on the graph is the leading order SM prediction for four top production of about 4 fb.

in the case of supersymmetric theories, and tW resonances for the theories with additional fermions. A systematic study of the characteristics of four top signals would be interesting, but is beyond the scope of this work.

5 Outlook

The top quark is a mystery. Its large mass may be an indication that is special in some way, and its status as the most recently discovered ingredient within the Standard Model allows great opportunity for new physics to cause its properties to deviate from our expectations. The truth about top is still largely unknown. In this article, we explore some of the deviations possible at the Tevatron and LHC when the top is not standard.

We find that bounds from top pair production at the Tevatron provide some information, but are relatively weak because of the possibility of cancellations between contributions from new physics. Kinematic distributions are much more powerful. The invariant mass distribution of the $t\bar{t}$ system can be pushed to peak at higher values, and the rapidity distribution may become more central or shifted asymmetrically, leading to a deviation in the forward-backward charge asymmetry of top pairs.

At the LHC, a more promising channel is provided by the production of four top quarks. The Standard Model predicts a very small rate for this process, leaving a lot of room for new physics to creep in. We find that a strongly coupled four top operator can lead to observable deviations even when its characteristic energy scale is multi-TeV.

While we are motivated by the picture that the top quark is composite, our results are presented in the language of effective field theory, and can be used far beyond the scenario at hand. *Any* theory of new physics modifying the inter-top interactions or top coupling to gluons is described at low energies by these operators. Once a deviation is found, the pattern of couplings reflects the UV physics which gave rise to the operators and provides the blue-print through which we can build a UV-complete description.

Acknowledgments

The authors are pleased to acknowledge conversations with J. Alwall and D. Choudhury. T Tait is grateful to the SLAC theory group for their excellent hospitality and over-all awesomeness while this work was being completed. Research at Argonne National Laboratory is supported in part by the Department of Energy under contract DE-AC02-06CH11357.

References

- [1] H. Georgi, L. Kaplan, D. Morin and A. Schenk, *Effects of top compositeness*, *Phys. Rev. D* **51** (1995) 3888 [[hep-ph/9410307](#)] [[SPIRES](#)].
- [2] B. Lillie, J. Shu and T.M.P. Tait, *Top compositeness at the Tevatron and LHC*, *JHEP* **04** (2008) 087 [[arXiv:0712.3057](#)] [[SPIRES](#)].
- [3] A. Pomarol and J. Serra, *Top quark compositeness: feasibility and implications*, *Phys. Rev. D* **78** (2008) 074026 [[arXiv:0806.3247](#)] [[SPIRES](#)].
- [4] M. Suzuki, *Model of a composite right-handed t quark: an alternative to the $t\bar{t}$ condensate model*, *Phys. Rev. D* **44** (1991) 3628 [[SPIRES](#)];
M.J. Strassler, *Generating a fermion mass hierarchy in a composite supersymmetric standard model*, *Phys. Lett. B* **376** (1996) 119 [[hep-ph/9510342](#)] [[SPIRES](#)];
A.E. Nelson and M.J. Strassler, *A realistic supersymmetric model with composite quarks*, *Phys. Rev. D* **56** (1997) 4226 [[hep-ph/9607362](#)] [[SPIRES](#)];
N. Arkani-Hamed, M. Porrati and L. Randall, *Holography and phenomenology*, *JHEP* **08** (2001) 017 [[hep-th/0012148](#)] [[SPIRES](#)];
T. Gherghetta and A. Pomarol, *The standard model partly supersymmetric*, *Phys. Rev. D* **67** (2003) 085018 [[hep-ph/0302001](#)] [[SPIRES](#)];
K. Agashe, A. Delgado, M.J. May and R. Sundrum, *RS1, custodial isospin and precision tests*, *JHEP* **08** (2003) 050 [[hep-ph/0308036](#)] [[SPIRES](#)];
A. Delgado and T.M.P. Tait, *A fat Higgs with a fat top*, *JHEP* **07** (2005) 023 [[hep-ph/0504224](#)] [[SPIRES](#)].
- [5] K. Agashe, A. Belyaev, T. Krupovnickas, G. Perez and J. Virzi, *LHC signals from warped extra dimensions*, *Phys. Rev. D* **77** (2008) 015003 [[hep-ph/0612015](#)] [[SPIRES](#)];
B. Lillie, L. Randall and L.-T. Wang, *The bulk RS KK-gluon at the LHC*, *JHEP* **09** (2007) 074 [[hep-ph/0701166](#)] [[SPIRES](#)];
B. Lillie, J. Shu and T.M.P. Tait, *Kaluza-Klein gluons as a diagnostic of warped models*, *Phys. Rev. D* **76** (2007) 115016 [[arXiv:0706.3960](#)] [[SPIRES](#)];
K. Agashe et al., *LHC signals for warped electroweak neutral gauge bosons*, *Phys. Rev. D* **76** (2007) 115015 [[arXiv:0709.0007](#)] [[SPIRES](#)];
K. Agashe, S. Gopalakrishna, T. Han, G.-Y. Huang and A. Soni, *LHC signals for warped electroweak charged gauge bosons*, [arXiv:0810.1497](#) [[SPIRES](#)].

- [6] A. Manohar and H. Georgi, *Chiral quarks and the nonrelativistic quark model*, *Nucl. Phys. B* **234** (1984) 189 [SPIRES];
H. Georgi, *Generalized dimensional analysis*, *Phys. Lett. B* **298** (1993) 187 [hep-ph/9207278] [SPIRES];
A.G. Cohen, D.B. Kaplan and A.E. Nelson, *Counting 4π 's in strongly coupled supersymmetry*, *Phys. Lett. B* **412** (1997) 301 [hep-ph/9706275] [SPIRES].
- [7] E. Eichten, K.D. Lane and M.E. Peskin, *New tests for quark and lepton substructure*, *Phys. Rev. Lett.* **50** (1983) 811 [SPIRES].
- [8] W. Buchmüller and D. Wyler, *Effective lagrangian analysis of new interactions and flavor conservation*, *Nucl. Phys. B* **268** (1986) 621 [SPIRES].
- [9] C.T. Hill and S.J. Parke, *Top production: sensitivity to new physics*, *Phys. Rev. D* **49** (1994) 4454 [hep-ph/9312324] [SPIRES].
- [10] D. Atwood, A. Kagan and T.G. Rizzo, *Constraining anomalous top quark couplings at the Tevatron*, *Phys. Rev. D* **52** (1995) 6264 [hep-ph/9407408] [SPIRES].
- [11] CDF collaboration, *Combination of top quark pair measurements cross section measurements*, CDF Note 8148, <http://www-cdf.fnal.gov/physics/new/top/top.html>.
- [12] D0 collaboration, V.M. Abazov et al., *Measurement of the $t\bar{t}$ production cross section in $p\bar{p}$ collisions at $\sqrt{s} = 1.96$ TeV*, *Phys. Rev. Lett.* **100** (2008) 192004 [arXiv:0803.2779] [SPIRES].
- [13] N. Kidonakis and R. Vogt, *Top quark production at the Tevatron at NNLO*, *Eur. Phys. J. C* **33** (2004) S466 [hep-ph/0309045] [SPIRES];
M. Cacciari, S. Frixione, M.L. Mangano, P. Nason and G. Ridolfi, *QCD analysis of first b cross-section data at 1.96 TeV*, *JHEP* **07** (2004) 033 [hep-ph/0312132] [SPIRES].
- [14] J. Alwall et al., *MadGraph/MadEvent v4: the new web generation*, *JHEP* **09** (2007) 028 [arXiv:0706.2334] [SPIRES].
- [15] R. Frederix and F. Maltoni, *Top pair invariant mass distribution: a window on new physics*, *JHEP* **01** (2009) 047 [arXiv:0712.2355] [SPIRES].
- [16] D0 collaboration, V.M. Abazov et al., *Search for narrow $t\bar{t}$ resonances in $p\bar{p}$ collisions at $\sqrt{s} = 1.8$ TeV*, *Phys. Rev. Lett.* **92** (2004) 221801 [hep-ex/0307079] [SPIRES];
CDF collaboration, T. Aaltonen et al., *Search for resonant $t\bar{t}$ production in $p\bar{p}$ collisions at $\sqrt{s} = 1.96$ -TeV*, *Phys. Rev. Lett.* **100** (2008) 231801 [arXiv:0709.0705] [SPIRES]; *Limits on the production of narrow $t\bar{t}$ resonances in $p\bar{p}$ collisions at $\sqrt{s} = 1.96$ TeV*, *Phys. Rev. D* **77** (2008) 051102 [arXiv:0710.5335] [SPIRES];
D0 collaboration, V.M. Abazov et al., *Search for $t\bar{t}$ resonances in the lepton plus jets final state in $p\bar{p}$ collisions at $\sqrt{s} = 1.96$ TeV*, *Phys. Lett. B* **668** (2008) 98 [arXiv:0804.3664] [SPIRES].
- [17] D0 collaboration, V.M. Abazov et al., *First measurement of the forward-backward charge asymmetry in top quark pair production*, *Phys. Rev. Lett.* **100** (2008) 142002 [arXiv:0712.0851] [SPIRES];
CDF collaboration, T. Aaltonen et al., *Forward-backward asymmetry in top quark production in $p\bar{p}$ collisions at $\sqrt{s} = 1.96$ TeV*, *Phys. Rev. Lett.* **101** (2008) 202001 [arXiv:0806.2472] [SPIRES].
- [18] M. Gerbush, T.J. Khoo, D.J. Phalen, A. Pierce and D. Tucker-Smith, *Color-octet scalars at the LHC*, *Phys. Rev. D* **77** (2008) 095003 [arXiv:0710.3133] [SPIRES].

- [19] D.E. Kaplan, K. Rehermann, M.D. Schwartz and B. Tweedie, *Top tagging: a method for identifying boosted hadronically decaying top quarks*, *Phys. Rev. Lett.* **101** (2008) 142001 [[arXiv:0806.0848](#)] [[SPIRES](#)];
 L.G. Almeida et al., *Substructure of high- p_T jets at the LHC*, [arXiv:0807.0234](#) [[SPIRES](#)];
 M. Cacciari, J. Rojo, G.P. Salam and G. Soyez, *Quantifying the performance of jet definitions for kinematic reconstruction at the LHC*, *JHEP* **12** (2008) 032 [[arXiv:0810.1304](#)] [[SPIRES](#)];
 L.G. Almeida, S.J. Lee, G. Perez, I. Sung and J. Virzi, *Top jets at the LHC*, [arXiv:0810.0934](#) [[SPIRES](#)];
 J. Thaler and L.-T. Wang, *Strategies to identify boosted tops*, *JHEP* **07** (2008) 092 [[arXiv:0806.0023](#)] [[SPIRES](#)].
- [20] B.A. Dobrescu, K. Kong and R. Mahbubani, *Massive color-octet bosons and pairs of resonances at hadron colliders*, *Phys. Lett. B* **670** (2008) 119 [[arXiv:0709.2378](#)] [[SPIRES](#)];
 C. Kilic, S. Schumann and M. Son, *Searching for multijet resonances at the LHC*, [arXiv:0810.5542](#) [[SPIRES](#)];
 C. Kilic, T. Okui and R. Sundrum, *Colored resonances at the Tevatron: phenomenology and discovery potential in multijets*, *JHEP* **07** (2008) 038 [[arXiv:0802.2568](#)] [[SPIRES](#)].
- [21] M. Guchait, F. Mahmoudi and K. Sridhar, *Associated production of a Kaluza-Klein excitation of a gluon with a $t\bar{t}$ pair at the LHC*, *Phys. Lett. B* **666** (2008) 347 [[arXiv:0710.2234](#)] [[SPIRES](#)].
- [22] R.N. Mohapatra, N. Okada and H.-B. Yu, *Diquark Higgs at LHC*, *Phys. Rev. D* **77** (2008) 011701 [[arXiv:0709.1486](#)] [[SPIRES](#)];
 C.-R. Chen, W. Klemm, V. Rentala and K. Wang, *Color sextet scalars at the CERN Large Hadron Collider*, *Phys. Rev. D* **79** (2009) 054002 [[arXiv:0811.2105](#)] [[SPIRES](#)].
- [23] J. Hisano, K. Kawagoe, R. Kitano and M.M. Nojiri, *Scenery from the top: study of the third generation squarks at CERN LHC*, *Phys. Rev. D* **66** (2002) 115004 [[hep-ph/0204078](#)] [[SPIRES](#)].
- [24] D. Choudhury, T.M.P. Tait and C.E.M. Wagner, *Beautiful mirrors and precision electroweak data*, *Phys. Rev. D* **65** (2002) 053002 [[hep-ph/0109097](#)] [[SPIRES](#)];
 C. Dennis, M. Karagoz Unel, G. Servant and J. Tseng, *Multi- W events at LHC from a warped extra dimension with custodial symmetry*, [hep-ph/0701158](#) [[SPIRES](#)];
 G.D. Kribs, T. Plehn, M. Spannowsky and T.M.P. Tait, *Four generations and Higgs physics*, *Phys. Rev. D* **76** (2007) 075016 [[arXiv:0706.3718](#)] [[SPIRES](#)];
 M. Carena, A.D. Medina, B. Panes, N.R. Shah and C.E.M. Wagner, *Collider phenomenology of gauge-Higgs unification scenarios in warped extra dimensions*, *Phys. Rev. D* **77** (2008) 076003 [[arXiv:0712.0095](#)] [[SPIRES](#)];
 R. Contino and G. Servant, *Discovering the top partners at the LHC using same-sign dilepton final states*, *JHEP* **06** (2008) 026 [[arXiv:0801.1679](#)] [[SPIRES](#)].